

# Double Hölder regularity of the hydrodynamic pressure in bounded domains

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[1] **L. De Rosa**, **M. Latocca** and **GS**, On double Hölder regularity of the hydrodynamic pressure in bounded domains, *Calc. Var. PDEs* 62 (2023), no. 3, Paper No. 85, 31.

[2] **L. De Rosa**, **M. Latocca** and **GS**, Full double Hölder regularity of the pressure in bounded domains (2023), submitted, preprint available at [arXiv:2301.06482](https://arxiv.org/abs/2301.06482).

## Euler equations and the hydrodynamic pressure

The evolution of an incompressible inviscid fluid is described by the **Euler equations**

$$\begin{cases} \partial_t u + \operatorname{div}(u \otimes u) + \nabla p = 0 & \text{in } \Omega \times (0, T), \\ \operatorname{div} u = 0 & \text{in } \Omega \times (0, T), \\ u \cdot n = 0 & \text{on } \partial\Omega \times (0, T), \end{cases} \quad (\text{E})$$

where:

- $\Omega \subset \mathbb{R}^d$ ,  $d \geq 2$ , is a bounded and simply connected  $C^2$  domain;
- $u: \Omega \times (0, T) \rightarrow \mathbb{R}^d$  is the **velocity**;
- $p: \Omega \times (0, T) \rightarrow \mathbb{R}$  is the **hydrodynamic pressure**;
- $n: \partial\Omega \rightarrow \mathbb{R}^d$  is the **outward unit normal** to  $\partial\Omega$ ;
- $u(\cdot, t) \cdot n = 0$  on  $\partial\Omega$  is the usual **no-flow condition**.

Hence, at each time, the pressure  $p$  solves the **elliptic Neumann problem**

$$\begin{cases} -\Delta p = \operatorname{div} \operatorname{div}(u \otimes u) & \text{in } \Omega, \\ \partial_n p = u \otimes u : \nabla n & \text{on } \partial\Omega, \end{cases} \quad (\text{P})$$

noticing that

$$\partial_n p = \nabla p \cdot n = -\operatorname{div}(u \otimes u) \cdot n = -u \cdot \nabla(u \cdot n) + u \otimes u : \nabla n = u \otimes u : \nabla n.$$

## Regularity problem for the pressure

Assume  $u \in C^{0,\gamma}(\Omega)$  with  $\operatorname{div} u = 0$  and  $u \cdot n = 0$  on  $\partial\Omega$  and consider

$$\begin{cases} -\Delta p = \operatorname{div} \operatorname{div} (u \otimes u) & \text{in } \Omega, \\ \partial_n p = u \otimes u : \nabla n & \text{on } \partial\Omega. \end{cases} \quad (\text{P})$$

By standard Schauder's estimates  $p \in C_{\text{loc}}^{0,\gamma}(\Omega)$ , but can we do better?

The regularity of the pressure **doubles** the regularity of the velocity!

**Theorem (Silvestre, 2011; Isett, 2023; Constantin, 2014; Colombo-De Rosa, 2020)**

$$\text{If } \Omega = \mathbb{R}^d \text{ or } \mathbb{T}^d, \text{ then } u \in C^{0,\gamma} \implies p \in \begin{cases} C^{0,2\gamma} & \text{if } 0 < \gamma < \frac{1}{2}, \\ \operatorname{Lip}_{\log} & \text{if } \gamma = \frac{1}{2}, \\ C^{1,2\gamma-1} & \text{if } \frac{1}{2} < \gamma < 1. \end{cases}$$

**Notation:**  $p \in \operatorname{Lip}_{\log}$  means  $|p(x) - p(y)| \leq C|x - y| \min\{1, |\log|x - y||\}$ .

**Remark:** See [Colombo-De Rosa-Forcella, 2020] for **Sobolev** and **Besov** spaces.

**Applications:** smoothness of trajectories of Euler flows, intermittency phenomena in turbulent flows, anomalous dissipation...

## Idea of proof in $\mathbb{R}^d$ (after Colombo-De Rosa)

The solution of  $-\Delta p = \operatorname{div} \operatorname{div} M$  decaying at infinity can be written as

$$p(x) = \int_{B_R(\bar{x})} \partial_{ij}^2 G_{\mathbb{R}^d}^D(x-y)(M^{ij}(y) - M^{ij}(x)) dy \\ - M^{ij}(x) \int_{\partial B_R(\bar{x})} \partial_i G_{\mathbb{R}^d}^D(x-y) \nu_j(y) dy$$

via **Green-Dirichlet kernel**  $G_{\mathbb{R}^d}^D$ , where  $M \in C_c^{0,\gamma}(\mathbb{R}^d; \mathbb{R}^{d \times d})$  with  $\operatorname{supp} M \subset B_R(x_0)$ .

Fix  $x_1, x_2 \in \mathbb{R}^d$  and  $\bar{x} = \frac{x_1+x_2}{2}$  and note that

$$\operatorname{div} \operatorname{div} (u \otimes u) = (\partial_i u^i)(\partial_j u^j) = \partial_i (u^i - u^i(x_1)) \partial_j (u^j - u^j(x_2)) \\ = \partial_{ij}^2 ((u^i - u^i(x_1))(u^j - u^j(x_2))),$$

so that we can write (for  $\gamma > \frac{1}{2}$  one takes derivatives  $\partial_k p$  and  $\partial_{ijk}^3 G_{\mathbb{R}^d}^D$ )

$$p(x_1) - p(x_2) \\ = \int_{B_R(\bar{x})} [\partial_{ij}^2 G_{\mathbb{R}^d}^D(x_1-y) - \partial_{ij}^2 G_{\mathbb{R}^d}^D(x_2-y)] (u^i(y) - u^i(x_1))(u^j(y) - u^j(x_2)) dy$$

and the conclusion follows by the decay properties of (derivatives of)  $G_{\mathbb{R}^d}^D$ .

## Pressure regularity on bounded domains

In the presence of a boundary, even the **single regularity** of the pressure is hard!

**Theorem (Bardos-Titi, 2022; Bardos-Boudros-Titi, 2023)**

$$d = 2 \text{ and } \partial\Omega \in C^2 \quad \text{or} \quad d = 3 \text{ and } \partial\Omega \in C^3 \implies p \in C^{0,\gamma}(\Omega)$$

Question: Can we still **double** the regularity on bounded sets?

**Theorem #1 (De Rosa-Latocca-S, 2023)**

$$d \geq 2 \text{ and } \partial\Omega \in C^{2,\delta} \implies p \in \begin{cases} C^{0,\gamma} & \text{if } 0 < \gamma \leq \frac{1}{2}, \\ C^{1,\min\{\delta, 2\gamma-1\}} & \text{if } \frac{1}{2} < \gamma < 1. \end{cases}$$

By interpolation, if  $\gamma < \frac{1}{2}$  then  $p \in C^{0,2\gamma-\varepsilon}$  for  $\varepsilon \in (0, 2\gamma)$  with  $d \geq 2$  and  $\partial\Omega \in C^{3,\delta}$ .

**Theorem #2 (De Rosa-Latocca-S, 2023)**

$$d \geq 2 \text{ and } \partial\Omega \in C^{2,1} \implies p \in \begin{cases} C^{0,2\gamma} & \text{if } 0 < \gamma < \frac{1}{2}, \\ C_*^1 & \text{if } \gamma = \frac{1}{2}. \end{cases}$$

Notation:  $C_*^1 \subsetneq \text{Lip}_{\log}$  with  $\|p\|_{C_*^1} = \|p\|_{L^\infty} + \sup_{x,h \in \mathbb{R}^d, h \neq 0} \frac{|p(x+h) + p(x-h) - 2p(x)|}{|h|}$ .

## Strategy of the proof of Theorem #1

Assume  $\Omega \subset \mathbb{R}^d$  is a simply connected bounded  $C^{2,\delta}$  domain.

### Approximation Lemma

If  $u \in C^{0,\gamma}(\Omega)$  satisfies  $\operatorname{div} u = 0$  and  $u \cdot n = 0$  on  $\partial\Omega$ , then there exist

$$u^\varepsilon \in C^\infty(\Omega) \cap C^{1,\delta}(\Omega) \quad \text{with} \quad \operatorname{div} u^\varepsilon = 0 \quad \text{and} \quad u^\varepsilon \cdot n = 0 \quad \text{on} \quad \partial\Omega$$

such that  $u^\varepsilon \rightarrow u$  in  $C^0(\bar{\Omega})$  as  $\varepsilon \rightarrow 0^+$  and

$$\sup_{\varepsilon > 0} \|u^\varepsilon\|_{C^{0,\gamma}(\Omega)} \leq C \|u\|_{C^{0,\gamma}(\Omega)}.$$

### Representation formula

Let  $u \in C^\infty(\Omega) \cap C^1(\bar{\Omega})$  be such that  $\operatorname{div} u = 0$  and  $u \cdot n|_{\partial\Omega} = 0$ . If  $p$  is a weak solution of (P), then

$$p(x) - \frac{1}{|\Omega|} \int_{\Omega} p(y) dy = \int_{\Omega} \partial_{y_i y_j} G_{\Omega}^N(x, y) (u_i(y) - u_i(x)) u_j(y) dy$$

for all  $x \in \Omega$ , where  $G_{\Omega}^N = G_{\Omega}^N(x, y)$  is the **Green-Neumann function** on  $\Omega$ .

We argue as in  $\mathbb{R}^d$ , after proving good **decay properties** of (derivatives of)  $G_{\Omega}^N$ .

## Strategy of the proof of Theorem #2 [1/3]

Now  $\gamma \in (0, \frac{1}{2}]$  and  $\Omega \subset \mathbb{R}^d$  be a simply connected bounded  $C^{2,1}$  domain.

By the **Approximation Lemma**, we may assume  $u \in C^\infty \cap C^{1,\delta}$ .

### Smoothed Theorem #2

If  $u \in C^\infty(\Omega) \cap C^{1,\delta}(\Omega)$  satisfies  $\operatorname{div} u = 0$  and  $u \cdot n|_{\partial\Omega} = 0$  on  $\partial\Omega$ , then **there exists a unique zero-average solution**  $p \in C^{1,\delta}(\Omega)$  of (P) such that

$$\|p\|_{C_*^{2\gamma}(\Omega)} \leq C_{\gamma,\Omega} \left( \|u\|_{C^{0,\gamma}(\Omega)}^2 + \|p\|_{L^\infty(\Omega)} \right).$$

Here  $C_*^{2\gamma} = C^{0,2\gamma}$  for  $\gamma < \frac{1}{2}$ . No big news for interior, just focus on **boundary!**

### Estimate near the boundary

If  $x_0 \in \partial\Omega$ , then exists a ball  $B_{R_0}(x_0)$  such that, with the notation above,

$$\|p\|_{C_*^{2\gamma}(\Omega \cap B_{R_0}(x_0))} \leq C_{\gamma,\Omega,R_0} \left( \|u\|_{C^{0,\gamma}(\Omega)}^2 + \|p\|_{L^\infty(\Omega)} \right).$$

Up to paying a suitable  $a \in \mathbb{R}$ , one reduces to  $q = p - \psi$  with  $\psi \in C^{1,\min\{\delta,\gamma\}}(\Omega)$  and

$$\begin{cases} -\Delta q = \operatorname{div} \operatorname{div} (u \otimes u) + a & \text{in } \Omega, \\ \partial_n q = 0 & \text{on } \partial\Omega. \end{cases}$$

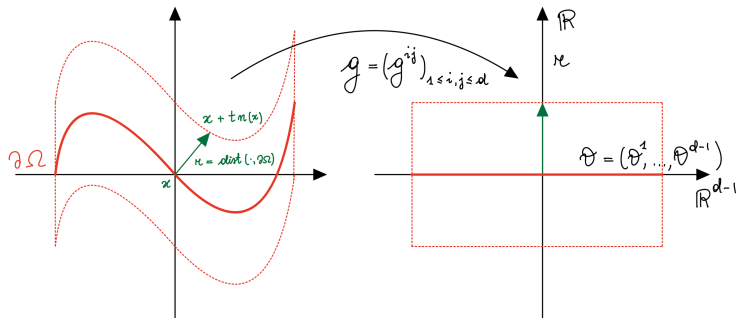
## Strategy of the proof of Theorem #2 [2/3]

### Stretching the boundary

We rewrite the problem in a local system of coordinates inside  $B_{R_0}(x_0)$  as

$$\begin{cases} -\partial_i (w g^{ij} \partial_j q) = \partial_{ij}^2 (w u^i u^j) + w a & \text{in } \Omega \cap U \\ \partial_r q(0, \vartheta) = 0 & \text{for } \vartheta \in \Theta. \end{cases}$$

where  $w = \sqrt{\det g(r, \vartheta)}$  and  $\partial_i$  are the derivatives in the normal coordinates  $r, \vartheta$ .





## Strategy of the proof of Theorem #2 [3/3]

We can reflect for  $r < 0$  and then extend to all  $\mathbb{R}^d$  with compact support.

The new extended solution  $\tilde{q}$  solves a problem of the form

$$-\chi \partial_i (\tilde{g}^{ij} \partial_j \tilde{q}) = \chi \partial_{ij}^2 (\tilde{w} \tilde{u}^i \tilde{u}^j) + \text{reminder} \quad \text{on } \mathbb{R}^d$$

where  $\chi$  is a suitable cut-off function.

The reminder term can be neglected, we just need to prove that

$$\|\tilde{q}\|_{C_*^{2\gamma}(\mathbb{R}^d)} \leq C \left( \|\tilde{w} \tilde{u}\|_{C^{0,\gamma}(\mathbb{R}^d)}^2 + \|\tilde{q}\|_{L^\infty(\mathbb{R}^d)} \right).$$

The idea is to study the **pseudodifferential operator**  $D_i = \tilde{g}^{ij} \partial_j$ . We may write

$$D = \underbrace{E}_{\text{elliptic}} + \underbrace{L}_{\text{lower order}}$$

and **invert** the elliptic part via Littlewood-Paley Analysis... painfully technical!

*Thank you for your attention!*

Slides available via [giorgio.stefani.math@gmail.com](mailto:giorgio.stefani.math@gmail.com) or [giorgiostefani.weebly.com](http://giorgiostefani.weebly.com).

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